

Static Viscoelastic Characteristics-Aggregation Structure Relationships  
of Monolayers at Air-Water Interface

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The melting temperature,  $T_m$  of fatty acid monolayers on the water surface was evaluated from the relationships between the subphase temperature dependence of static viscoelasticities and the aggregation structure of those monolayers.  $T_m$ s of myristic, palmitic, and stearic acids in a monolayer state are ca. 278 K, 301 K, and 317 K, respectively. These values are 26-50 K lower than those in respective three-dimensional crystalline states. With an increase of subphase temperature, the aggregation state of those monolayers on the water surface varied from an oriented crystalline state, an unoriented crystalline one, an amorphous one and finally, a dissoluble unstable state.

It has been generally accepted that phase transitions of gas, liquid, liquid condensed, and solid phases occur upon compressing fatty acid molecules on the water surface. However, it was recently proposed that two-dimensional domains were grown at the surface pressure of  $0 \text{ mN}\cdot\text{m}^{-1}$  right after spreading a solution of fatty acid on the water surface.<sup>1,2)</sup> The aggregation state of monolayer strongly depends on the value of spreading temperature,  $T_{sp}$ , in relative to the melting temperature,  $T_m$ , of the monolayer on the water surface. In this study, the temperature dependence of static viscoelastic properties of monolayers was investigated in conjunction with morphological observation and electron diffraction study of monolayers. The crystalline relaxation behavior and the melting behavior of monolayers on the water surface were also investigated.

Benzene solutions of myristic ( $C_{14}$ ), palmitic ( $C_{16}$ ) and stearic ( $C_{18}$ ) acids were prepared at the concentrations of  $3.7 \times 10^{-3}$ ,  $2.8 \times 10^{-3}$ , and  $2.2 \times 10^{-3} \text{ mol}\cdot\text{l}^{-1}$ , respectively. Surface pressure-area ( $\Pi$ -A) isotherms were measured at various  $T_{sp}$  with a Micro Proseccer Film Balance System (Sanesu Keisoku Co.Ltd.). The static elasticity,  $K_S$  of monolayer on the water surface was evaluated from the  $\Pi$ -A isotherm by using the following equation.

$$K_S = -A(d\Pi/dA)$$

Figure 1 shows the  $\Pi$ -A (solid line) and the  $\log K_S$ -A (broken line) isotherms for the stearic acid monolayer at  $T_{sp}$  of 293 K. The regions of a-b, b-c, c-d in the  $\log K_S$ -A isotherm correspond to the phases which are generally called a liquid expanded, a liquid condensed, and a solid phase, respectively. The maximum of  $\log K_S$ ,  $\log K_{S(\max)}$  was observed around the collapsing point (d), where molecules

in the monolayer were packed most closely, although locally collapsed domains overlapped the monolayer.<sup>1)</sup> In this paper, the temperature dependence of  $\log K_{S(\max)}$  was adopted to determine  $T_m$  of the monolayer as well as the characteristic viscoelastic behaviors of the monolayer. The monolayer was transferred onto collodion-covered copper grids at various  $T_{sp}$  by a vertical dipping method with the speed of  $40 \text{ mm}\cdot\text{min}^{-1}$ , at the surface pressures where each monolayer was morphologically homogeneous.<sup>1)</sup> Bright field electron micrographs and electron diffraction (ED) patterns were taken with a Hitachi H-500 electron microscope. Pt-carbon was vapor-deposited onto the monolayer samples with the shadowing angle of  $25^\circ$  for the bright field electron microscopic observation.

Figure 2 shows the  $T_{sp}$  dependence of  $\log K_{S(\max)}$  for the stearic acid monolayer on the water surface and the ED patterns of the monolayer transferred at the surface pressure of  $20 \text{ mN}\cdot\text{m}^{-1}$ . The slope of the  $\log K_{S(\max)}$  vs.  $T_{sp}$  curve fairly changed at ca. 298 K. In a  $T_{sp}$  region lower than 298 K, the ED pattern was a very sharp hexagonal spot, which indicated a regularly packed crystalline state. On the other hand, at a higher  $T_{sp}$  region, the electron diffraction spot tended to be arc along an azimuthal direction with an increase of  $T_{sp}$ , resulting in the Debye ring. This indicates that molecular and/or crystalline orientations become broader or less regular with an increase of  $T_{sp}$  above ca. 298 K. The value of  $\log K_{S(\max)}$  decreased abruptly at ca. 317 K and the ED pattern at 318 K showed an amorphous halo. This indicates that  $T_m$  of stearic acid monolayer on the water surface is ca. 317 K.

Figure 3 shows the  $T_{sp}$  dependence of  $(10\bar{1}0)$  spacing evaluated from ED

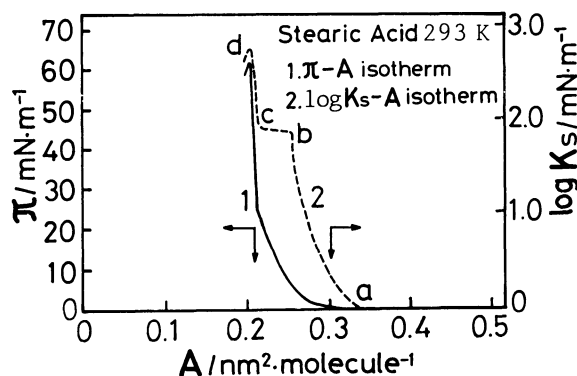
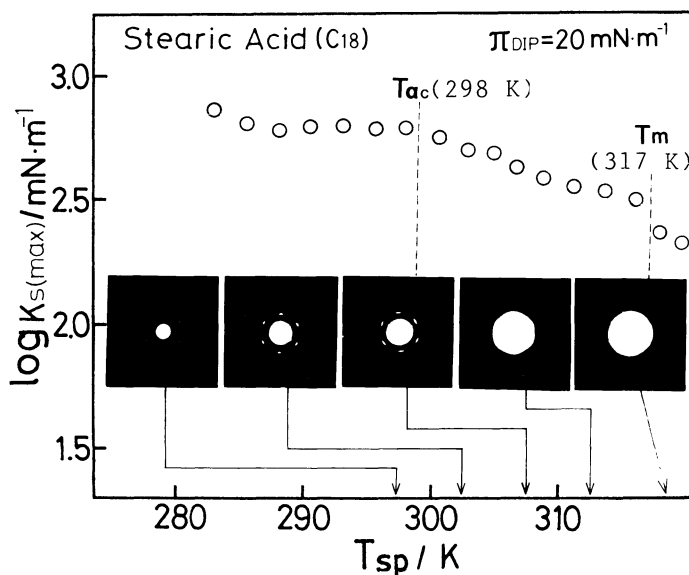


Fig. 1.  $\pi$ -A and  $\log K_S$ -A isotherms of stearic acid monolayer.

Fig. 2.  $T_{sp}$  dependence of  $\log K_{S(\max)}$  and ED patterns of stearic acid monolayers transferred at a surface pressure of  $20 \text{ mN}\cdot\text{m}^{-1}$ .



patterns for stearic acid monolayers of which crystal system was considered to be hexagonal. The  $(10\bar{1}0)$  spacing vs.  $T_{sp}$  curve exhibited a distinct break at ca. 298 K; the thermal expansion coefficient of crystal lattice became greater above this temperature. In general, the thermal expansion of the lattice spacing starts to deviate from an approximate linear relationship with measuring temperature when abrupt changes of physical or viscoelastic properties in the crystalline region cause an increase in contribution from the anharmonic term for the intermolecular potential energy.<sup>3)</sup> Therefore, Figs. 2 and 3 indicate that the thermal molecular motion in the stearic acid crystalline region must be in close relation to the viscoelastic crystalline relaxation process; this will be accompanied by the remarkable break of the thermal expansion of the  $(10\bar{1}0)$  spacing.

Figure 4 shows the  $T_{sp}$  dependence of  $\log K_S(\max)$  with the electron micrographs and ED patterns for the myristic acid monolayer transferred at  $15 \cdot \text{mN} \cdot \text{m}^{-1}$ . The value of  $\log K_S(\max)$  decreased abruptly at ca. 278 and 294 K. The ED patterns at 274 K and 288 K showed a crystalline ring and an amorphous halo, respectively. Therefore,  $T_m$  of myristic acid monolayer at the air-water interface is ca. 278 K. The ED pattern and the bright field electron micrograph of the monolayer at 303 K showed an amorphous halo and a extremely uniform structure, respectively. In a  $T_{sp}$  region higher than 294 K, the limiting area of the  $\pi$ -A isotherm became smaller

Fig. 4.  $T_{sp}$  dependence of  $\log K_S(\max)$ , electron micrographs, and ED patterns of myristic acid monolayers transferred at a surface pressure of  $15 \text{ mN} \cdot \text{m}^{-1}$ .

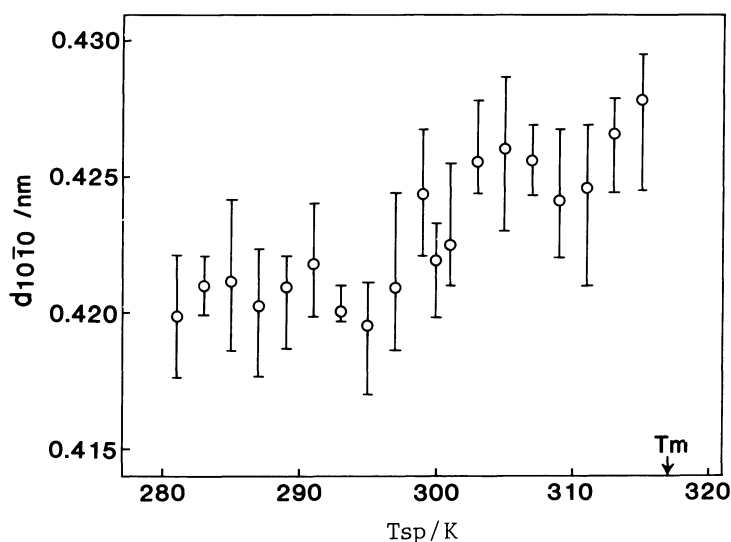
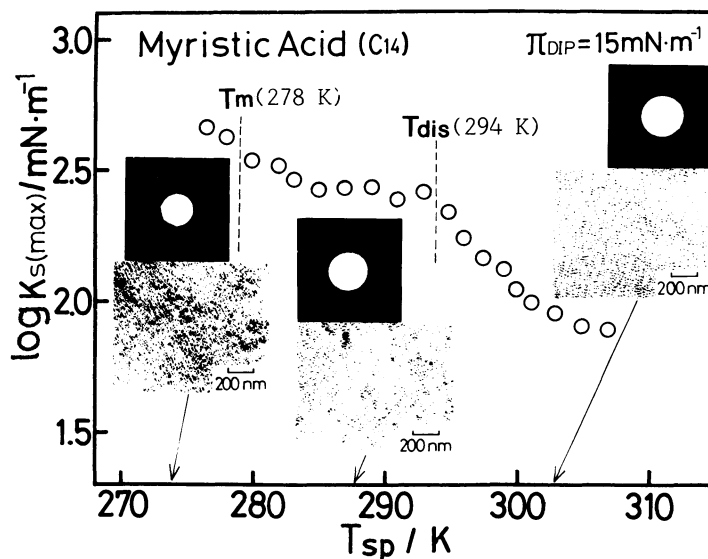


Fig. 3.  $T_{sp}$  dependence of  $(10\bar{1}0)$  spacing of stearic acid monolayers transferred at a surface pressure of  $20 \text{ mN} \cdot \text{m}^{-1}$ .



with an increase of  $T_{sp}$  or a standing time on the water surface after spreading a solution. This indicates a decrease in the number of myristic acid molecules on the water surface, probably owing to their dissolution into the water phase. Also, in the case of palmitic acid monolayer, the crystalline relaxation temperature,  $T_{\alpha c}$  (291 K),  $T_m$  (301 K) and  $T_{dis}$  (308 K) were evaluated on the basis of  $\log K_S(\max) - T_{sp}$  behavior.

Figure 5 shows the variations of  $T_m$  with alkyl chain length for fatty acids in a monolayer state, and a three-dimensional crystalline one based on the DSC measurement together

with those of n-paraffin in a three-dimensional crystalline one.  $T_m$  of n-paraffin was evaluated from the Broadhurst's equation.<sup>4)</sup>  $T_m$ s of myristic, palmitic, and stearic acids in a monolayer state were ca. 278 K, 301 K, and 317 K, respectively. It is difficult to evaluate the accurate  $T_m$  of fatty acids with alkyl chain longer than  $C_{18}$  in a monolayer state due to equipmental limitation on increasing  $T_{sp}$ .  $T_m$ s of fatty acids in a monolayer state are much lower than those in a three-dimensional crystalline one as shown in Fig.5. However, the difference between  $T_m$ s in these two states decreased with an increase of alkyl chain length; this indicates the formation of more stable monolayer due to an increase of intermolecular aggregation force with increasing alkyl chain length.

In conclusion, the  $T_{sp}$  dependence of  $\log K_S(\max)$  was strongly related to the phase transitions of monolayer or thermal motion of amphiphilic molecules.  $T_{\alpha c}$  and  $T_m$  of monolayer can be evaluated from the temperature dependence of  $\Pi$ -A isotherm and morphological or structural informations at different  $T_{sp}$ . It became clear that  $T_m$  of a monolayer at air-water interface was much lower than that in a three-dimensional crystalline state.

#### References

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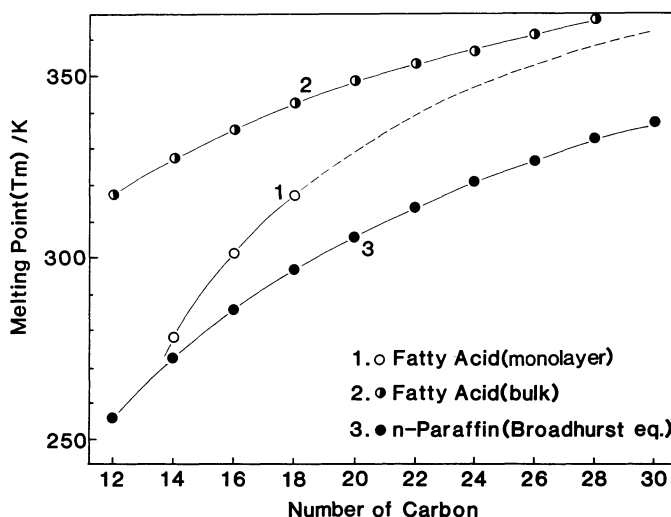


Fig. 5. Variation of  $T_m$  with alkyl chain length for (1) fatty acids in a monolayer state; (2) fatty acids in a three-dimensional crystalline state; and (3) n-paraffin in a three-dimensional crystalline state.

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